

INFLATION FROM AXION MONODROMY

E. PATER

INTRO inflation: $\ddot{\phi} > 0 \Leftrightarrow \frac{d}{dt} \left(\frac{\dot{\phi}}{\lambda_{Hubble}} \right) > 0$

• Many vanilla models compatible w/ WMAP. How to distinguish?

Exciting signatures: NON-GAUSSIANITY & TENSORS

$f_{NL} < \mathcal{O}(100)$ $r < 0.2$
1% accuracy upper bound

Lytic bound: $\frac{d\phi}{M_p} \sim dN \sqrt{2\epsilon} \sim dN \sqrt{\pi/8}$ (slow roll)

SuperPlanckian range $\Rightarrow \frac{\Delta\phi}{M_p} > \sqrt{\frac{\pi}{0.01}} \frac{N_{obs}}{30}$

EFT approach: study UV-sensitive observables.

Inflation is UV-sensitive: $V(\phi) = \frac{1}{2} m^2 \phi^2 + \sum_n \frac{c_n}{M_p^{n-4}} \phi^n$

- Large field is natural iff \exists symmetry, e.g. shift sym
- Small field requires tuning dim 5 & dim 6 operators.

Axions: $\mathcal{L}[\phi] = \mathcal{L}[\phi + c] \Rightarrow$ only $\partial\phi$ couplings

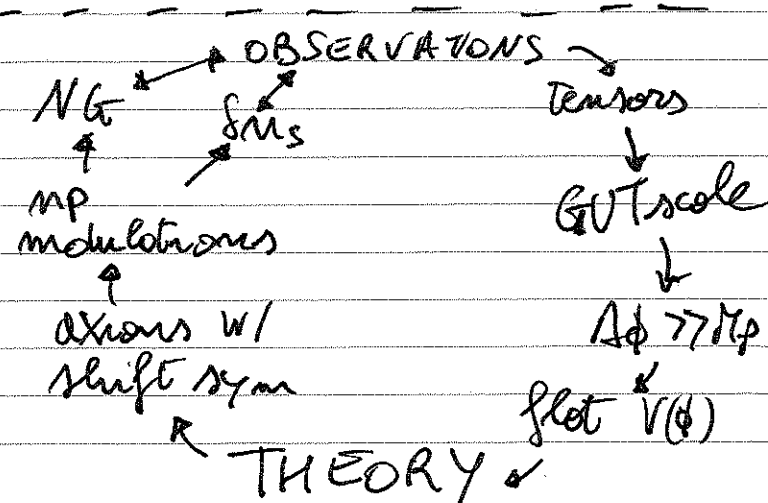
Beyond p-theory: $V(\phi) = \text{const} + \Lambda^4 \cos(\phi/f)$

$f \equiv$ AXION DECAY CONST ; $\Lambda \sim e^{-\gamma g}$ NON PERTURBATIVE

AXION MONODROMY

$$V(\phi) = V_{SR}(\phi) + \Lambda^4 \cos(\phi/f)$$

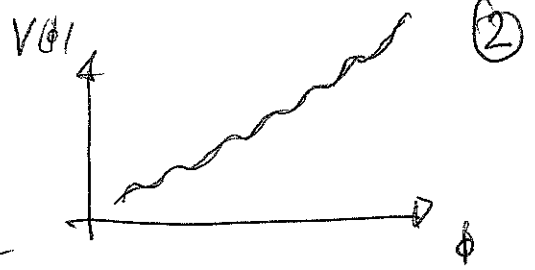
- Embedded in ST
- TeV scale: $r = 0.07$
- Peculiar: f_{NL} & f_2 modulated.



The Model

$$V = V_{SR}(\phi) + \Lambda^4 \cos(\phi/f)$$

R slow-roll flat potential



- Λ dimension M^4 . Non-perturbatively generated
- f axion decay constant. Dimension M^1 ,
In string theory $f \approx \frac{g_s^\#}{L^\#} M_p$ so always $f \ll M_p$

Trade for other two parameters.

- $b \equiv \frac{\Lambda^4}{V'_{SR} f}$. Then $V(\phi)$ is monotonic iff $b < 1$
- $\alpha \equiv \frac{\omega}{H} \equiv \frac{\sqrt{2\epsilon}}{f} = \frac{|\dot{\phi}|}{fH}$. ω is frequency of $\cos(\phi/f)$

While H is frequency of dS background.

$\alpha = \#$ of oscillations in CMB.

To have any interesting effect $d \gg 1$

Background

$$\phi(t) = \phi_0(t) + \phi_{osci} \sin\left(\frac{\phi_0(t)}{f}\right)$$

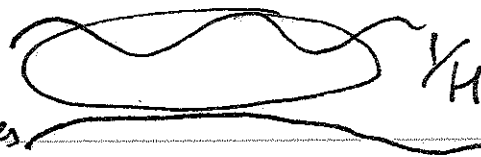
- $\phi_0(t)$ is the slow-roll background. $\mathcal{O}(b^0)$

- Due to $\phi_{osci} \sim \mathcal{O}(b)$ the dynamics violates slow roll.

$$\begin{cases} \epsilon = \epsilon_0 + \epsilon_{osci} \\ \eta \approx \eta_0 + \eta_{osci} \\ \dots \end{cases} \quad \text{Because} \quad \frac{1}{H} \frac{\partial}{\partial t} \sim \alpha \gg 1$$

$$\epsilon_{osci} \ll \eta_{osci} \ll \frac{\eta}{H} \ll \dots$$

$x \gg x_{res}$ x_{res} $x \ll x_{res}$



PERTURBATIONS

Comoving gauge $g_{ij} = \delta_{ij} a^2 (1 + 2R)$ $\delta\phi = 0$

M-S eq: $R_k'' - \frac{2(1+2\epsilon+\delta)}{x} R_k' + R_k = 0$ w/ $\kappa = -kz$

$\epsilon(x), \delta(x)$ oscillate very fast. Slow-roll is violated!

$(\epsilon \equiv -\frac{\dot{H}}{H^2}; \delta \equiv \frac{\ddot{H}}{2H\dot{H}})$ $E_{osc} \ll \delta_{osc} \sim -3b \sin(\phi/f)$

$\phi/f = \phi_k/f + \alpha \log x$

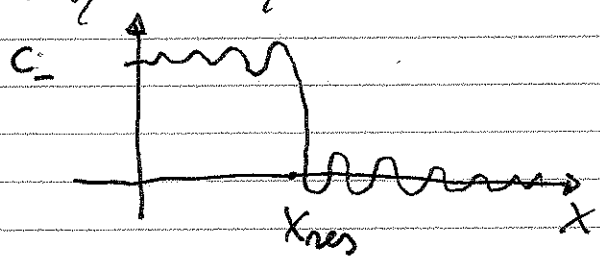
$R_k \sim e^{ix} \Rightarrow R_k$ & $\delta(x)$ RESONATE for $x = x_{res} = \alpha/2$

$x \gg x_{res}$ $R_k(x) = R_k^{(0)} \sqrt{\frac{\pi}{2}} x^\nu H_\nu^{(1)}(x) \approx R_k^{(0)} e^{ix} (1 - ix)$

$x \ll x_{res}$ $R_k(x) = R_k^{(0)} \sqrt{\frac{\pi}{2}} x^\nu [c_k^+ H_\nu^{(1)}(x) - c_k^- H_\nu^{(2)}(x)]$

Easy analytic solution.

$c_k \rightarrow 3b \sqrt{\frac{\pi}{2\alpha}} e^{-\alpha \phi_k/f}$

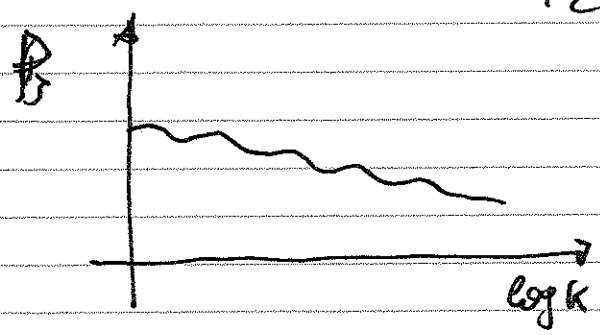


RESONANCE!

under horizon
↓
neglect gravity

POWER SPECTRUM

$\langle R_k^2 \rangle = (2\pi)^3 f^3 \frac{H^2}{4\epsilon} [1 + 3b \sqrt{\frac{2\pi}{\alpha}} \cos(\alpha \log \frac{k}{k_*} + \theta)]$



Period $\frac{2\pi}{\alpha} \Rightarrow \alpha$ in CMB scales

Amplitude $3b \sqrt{\frac{2\pi}{\alpha}}$
(small w/b, large w/ $\alpha \gg 1$)

Constraints from WMAP 5 $f_b < 10^{-4}$

INTERACTIONS

Interactions come from $\mathcal{L}[\phi] \Rightarrow$ can neglect gravity.

Decoupling (limit) $S = \int \sqrt{-g} [R + \mathcal{L}(\phi)]$

- Easier computations, eg. DBI, ghost, EFT of ...
- Compute uncomputable quantities, e.g. N -point.

BISPECTRUM

$\langle R(k_1) R(k_2) R(k_3) \rangle = 9-3-3$ dof.

Text of NG:

$\langle R^3 \rangle_c = \int_{\text{shape}} F(k_1, k_2, k_3) \cdot (2\pi)^3 \delta^3$

$H_{\pm} = \int d^3x \frac{v^{(n)}}{n!} \phi^n + \text{mixing w/ gravity}$

Explicit inspection shows that mixing is subleading.

$\langle R^3 \rangle = -i \int_{-\infty}^t dt' \langle [R^3, H_{\pm}] \rangle$ w/ $R_k = R_k^{(0)} x^{3/2} H_{3/2}^{(1)}$

Schematically: $\int dx e^{-\alpha x} \text{num}(\phi/k_f + d \log x)$

b/c $v^{(3)} \approx \frac{3\alpha b}{f} \text{num}(\phi/f) \rightarrow$

Analytic exact solution or stationary phase approx:

$\langle R^3 \rangle = \frac{(2\pi)^7 \Delta^4 \delta^3}{\pi_i k_i^2} f^{\text{res}} \left[\text{num}(d \log \frac{k_i}{k_j}) + \frac{K}{\alpha} \sum \frac{1}{k_i} \text{cos}(d \log) \right]$

$f^{\text{res}} = \frac{3b\sqrt{2\pi}}{f} \alpha^{3/2}$

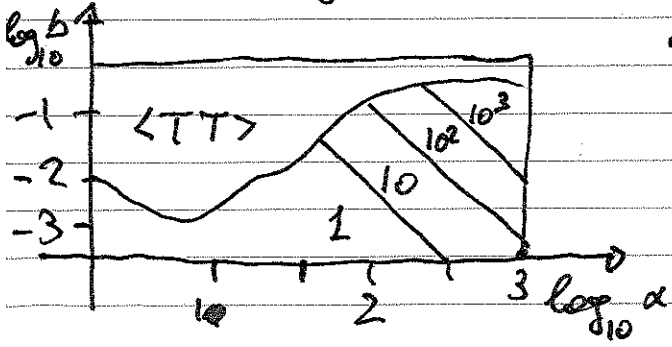
• linear in $b \ll 1$ & $\alpha^{3/2} \gg 1$

• $\delta M_s \sim b/\sqrt{\alpha}$ complementary

• $F^{\text{res}}(k_i)$ is not scale-invariant

• $F^{\text{res}} \perp$ only other knows shape

$\text{cos} < \frac{\pi}{\alpha} \ll 1$



(5)

N-POINT

EP. Z. KEBLOND
1010.4565

The simplest model! Computable!

$$\langle \delta\phi^N \rangle = \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \dots$$

at tree level

$$\langle \text{diagram 1} \rangle = (2\pi)^3 \frac{A_N}{K^{N-3} \prod_i k_i^2} \left[\text{Im} \left(\frac{\phi_K}{f} \right) - \frac{K}{d} \sum \frac{1}{k_i} \cos \left(\frac{\phi_i}{f} \right) \right]$$

$$A_N = (-)^N \frac{3b\sqrt{2\pi}}{2} \alpha^{2N-9/2} (2\pi^{2/3} \Delta_R^2)^{N-1}$$

- Linear in b & α^{2N}
- Multi-vertex gets $b\sqrt{\alpha} \sum_{n=1}^3 \left(\frac{K}{q\alpha} \right)^n$ for each vertex plus a combinatoric factor $C(N)$.
- Estimate $C(N) \Rightarrow$ Multi-v. suppressed to $N=15-20$
- Computed (first time) $\langle \delta\phi^N \rangle$ w/ $N \gg 3$. (e.g. $N=4$)
- Can check a large number of consistency relations.

OUTLOOK

- Inflation is UV sensitive. Large field even more.
- Axion monodromy provides a UV embedding in ST.
- $r=0.07$ & potentially f_{NL}^{res} . Testable.
- New NG, \perp to the others.
- The simplest model, $\langle \delta\phi^N \rangle$ w/ $N \gg 3$!